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“Small-Scale Dynamo Action in Primordial Halos”

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Göttingen 2012

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2. Our Model for the Primordial ISM

2.1 Magnetic Seed Fields

2.2 Our Approach to Turbulence

2.3 Model for the Primordial Gas

3. Magnetic Field Amplification by the Small-Scale Dynamo

3.1 Kinematic Growth

3.2 Non-Linear Growth

4. Conclusion

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Small-Scale Dynamo

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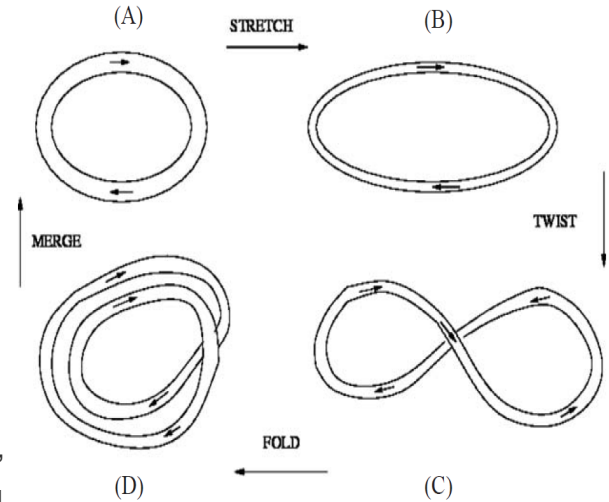
3. Small-Scale Dynamo

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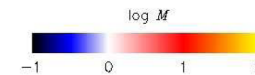
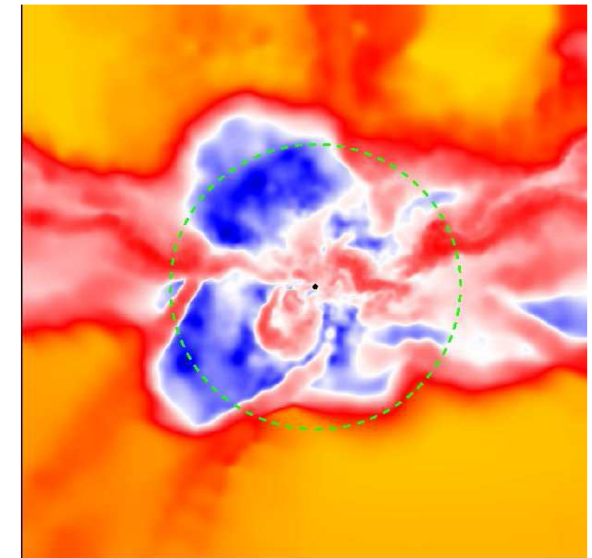
- **small-scale dynamo:** process which converts kinetic energy from turbulence into magnetic energy on short timescales

“Stretch-Twist-Fold Model”
[Brandenburg & Subramanian 2005]



- first point in time, where dynamo can operate:
formation of the first stars
(first turbulent halos)

Mach number in Primordial Halo
[Greif et al. 2008]



Size: 40 kpc (comoving)
x-y plane
z = 10.62
t_H = 429.4 Myr

Primordial Star Formation

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2. Primordial ISM

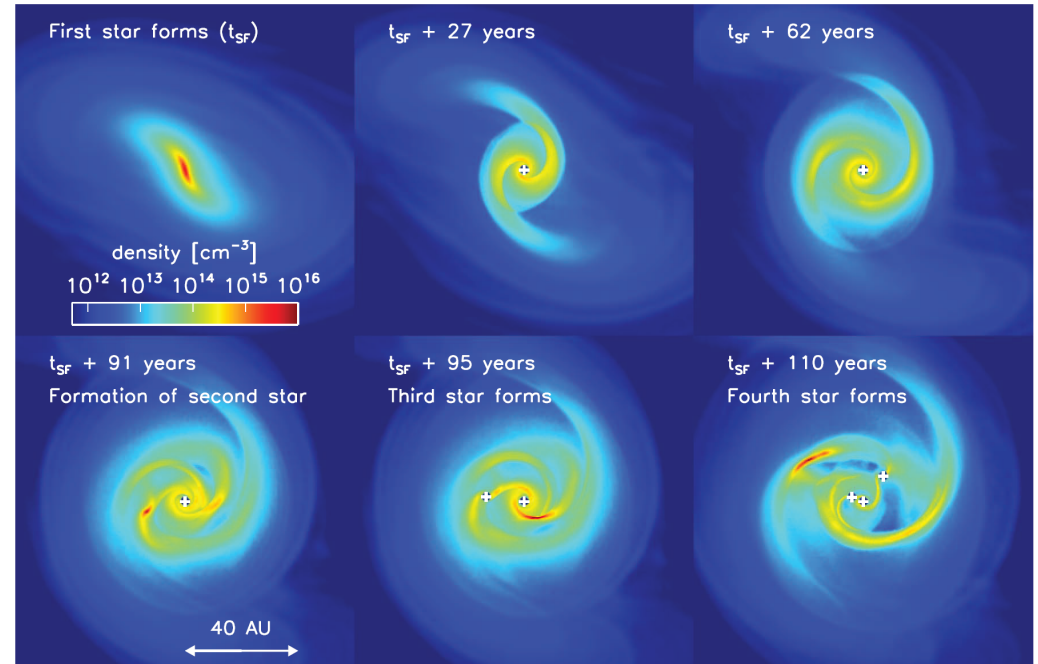
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- start at $z \approx 20$
- formation in dark-matter mini halos
- gas composition: H, He, Li, ...
- temperature:
 $T \approx 1000$ K



[Clark et al., Science, 2011]

- further ingredients:
 - weak magnetic seed fields
 - turbulence: mixture of solenoidal and compressive modes

=> dynamo action in principle possible

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Magnetic Seed Fields

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- **inflation:**

$$B \approx (10^{-34} - 10^{-9}) \text{ G} \quad [\text{e.g. Turner \& Widrow 1988}]$$

- **phase transitions in the early Universe:**

$$B \approx 10^{-20} \text{ G} \quad [\text{QCD phase transition, 10 Mpc comoving scale, e.g. Sigl et al. 1997}]$$

- **battery processes:**

$$B \approx 10^{-20} \text{ G} \quad [\text{Biermann battery, kpc scale, e.g. Xu et al. 2008}]$$

=> extremely weak seed fields!

Turbulence

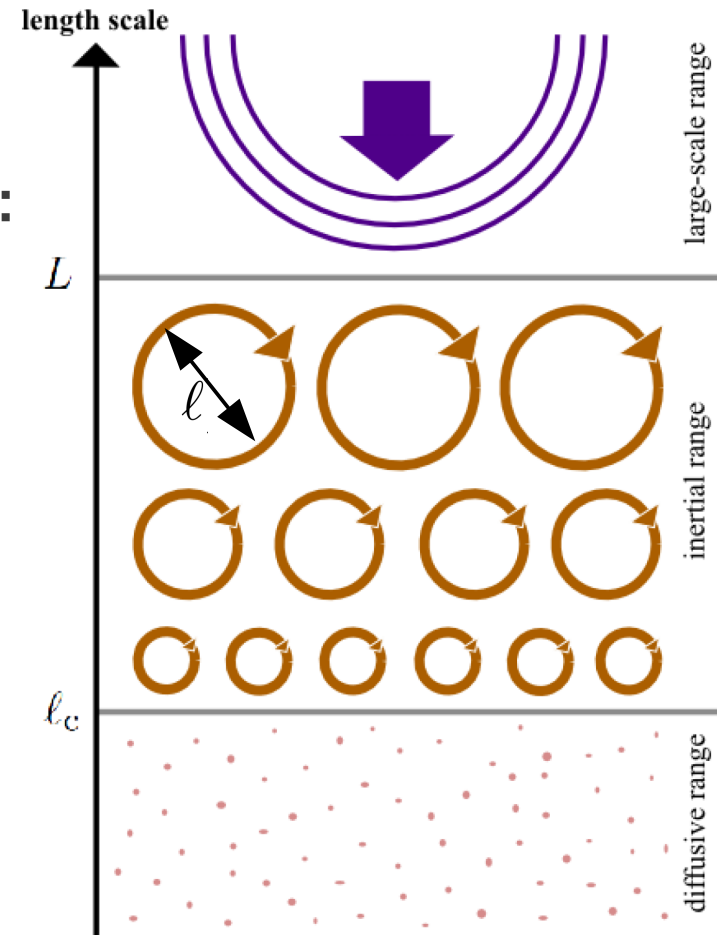
- turbulence in first star-forming halos is **driven by accretion**
- driving scale: Jeans scale
- **different types of turbulence:**
relation in the inertial range:

$$\delta v(l) \propto l^\vartheta$$

$$1/3 \leq \vartheta \leq 1/2$$

Kolmogorov
(incompressible)

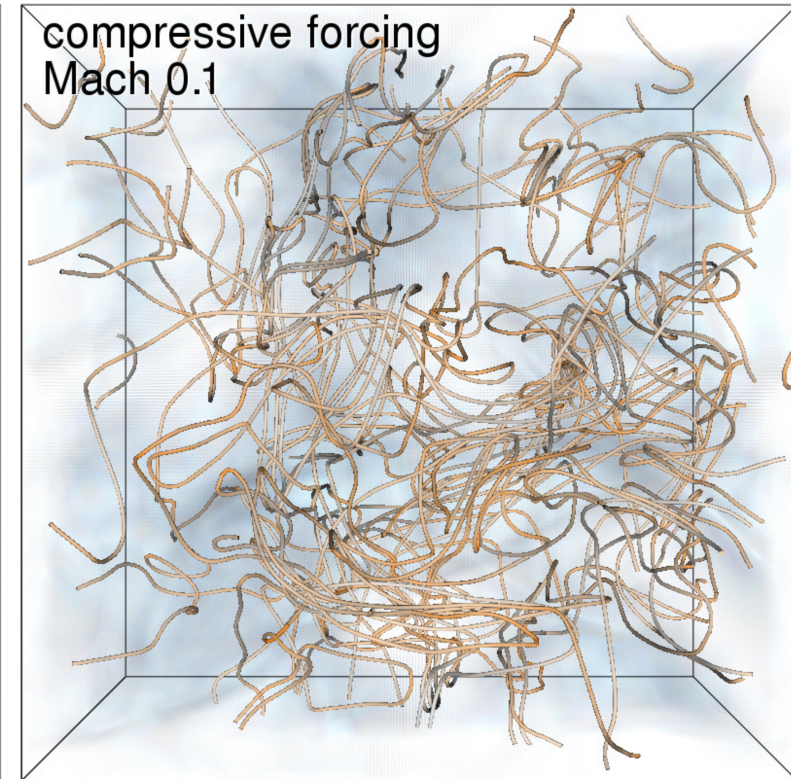
Burgers
(highly compressible)



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Different Forcing/Turbulence

effect of different types of forcing on the structure of magnetic field lines:



[Federrath et al., PRL, 2011]

=> solenoidal forcing twists field lines more effectively

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The Velocity Field

- **separation of smooth and turbulent component:**

$$\mathbf{v} = \langle \mathbf{v} \rangle + \delta \mathbf{v}$$

- **properties of turbulent field $\delta \mathbf{v}$:**
 - isotropic and homogeneous
 - Gaussian random field with zero mean
 - delta-correlated in time

- **spatial two-point correlation function**
= “measure for turbulent kinetic energy density”:

$$\begin{aligned} & \langle \delta v_i(\mathbf{r}_1, t) \delta v_j(\mathbf{r}_2, s) \rangle \\ &= \left(\left(\delta_{ij} - \frac{r_i r_j}{r^2} \right) T_N(r) + \frac{r_i r_j}{r^2} T_L(r) \right) \delta(t - s) \end{aligned}$$

\uparrow
transversal

\uparrow
longitudinal

(note: We neglect helicity.)

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Model for T_L and T_N

- motivation: diffusion coefficient in inertial range

- **model for general turbulence:**

$$T_L(r) = \begin{cases} \frac{VL}{3} \left(1 - Re^{(1-\vartheta)/(1+\vartheta)} \left(\frac{r}{L} \right)^2 \right) & 0 < r < \ell_c \\ \frac{VL}{3} \left(1 - \left(\frac{r}{L} \right)^{\vartheta+1} \right) & \ell_c < r < L \\ 0 & L < r \end{cases}$$

[Vainsthein, Sov. Phys. JETP, 1982; Subramanian, ArXiv, 1997]

$$T_N(r) = \begin{cases} \frac{VL}{3} \left(1 - t(\vartheta) Re^{(1-\vartheta)/(1+\vartheta)} \left(\frac{r}{L} \right)^2 \right) & 0 < r < \ell_c \\ \frac{VL}{3} \left(1 - t(\vartheta) \left(\frac{r}{L} \right)^{\vartheta+1} \right) & \ell_c < r < L \\ 0 & L < r \end{cases}$$

[Schober et al., PRE, 2012]

(ℓ_c : viscous scale, L : scale of largest fluctuations,

$Re = \frac{VL}{\nu}$: Reynolds number)

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Chemistry & MHD Quantities

- **one-zone model of Glover & Savin (2009)**
(with modification connecting collapse time and equation of state [Schleicher et al. 2009] and additional Li-chemistry [Bovino et al. 2011]):
~ 30 different species (atomic and molecular)
~ 400 different chemical reactions

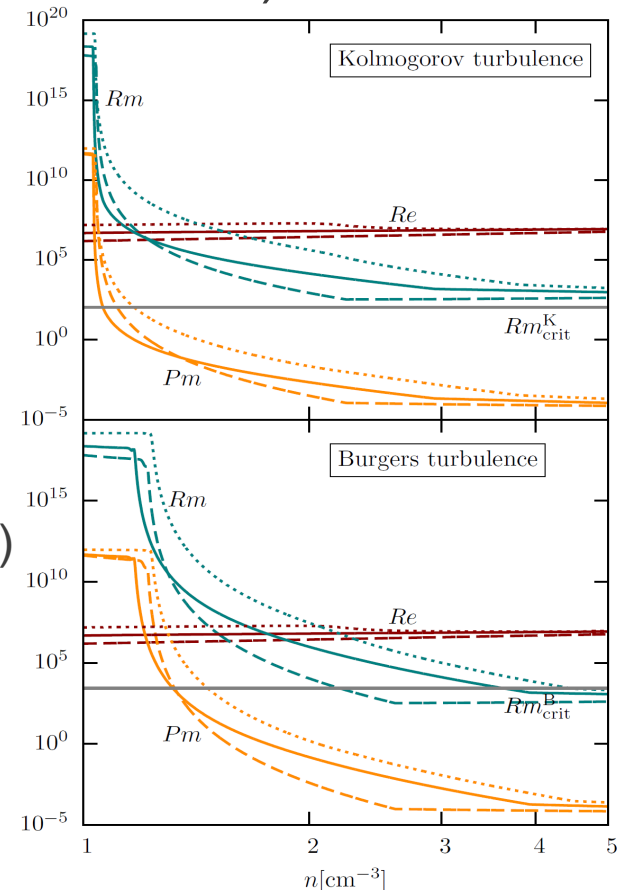
- **characteristic MHD quantities:**

$$Re = \frac{VL}{\nu} \quad \text{hydrodynamical Reynolds number}$$

$$Rm = \frac{VL}{\eta} \quad \text{magnetic Reynolds number} \\ (\eta = \eta_{\text{Ohm}} + \eta_{\text{AD}}(\mathbf{B}))$$

$$Pm = \frac{Rm}{Re} = \frac{\nu}{\eta} \quad \text{magnetic Prandtl number}$$

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3. Magnetic Field Amplification by the Small-Scale Dynamo

MHD-Dynamos – Mathematical Description

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- idea: divide magnetic field into **mean and turbulent component**

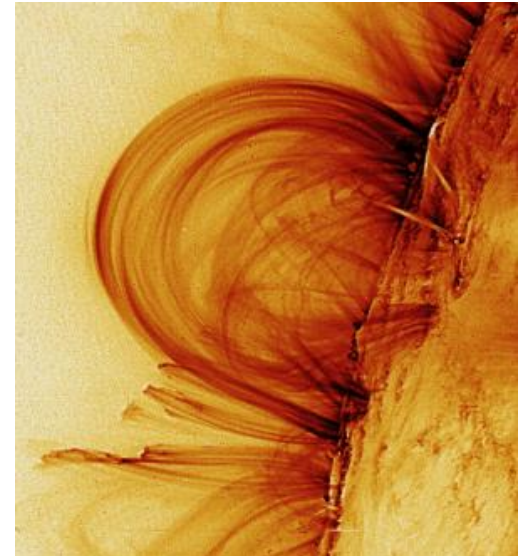
$$\mathbf{B} = \langle \mathbf{B} \rangle + \delta \mathbf{B}$$

- put into **induction equation**:

$$\frac{\partial \mathbf{B}}{\partial t} = \nabla \times \mathbf{v} \times \mathbf{B} - \eta \nabla \times \nabla \times \mathbf{B}$$

=> evolution equations for mean
and turbulent field
(large-scale dynamo and
small-scale dynamo)

Magnetic active region on the Sun
[credit: NASA, Trace Mission]



Kazantsev Theory: Model for Kinematic Growth

- put $\langle \delta B_i(\mathbf{r}_1, t) \delta B_j(\mathbf{r}_2, t) \rangle$ (= “magnetic energy density”) into induction equation

=> **Kazantsev equation** [Kazantsev, 1968]:

$$-\kappa_{\text{diff}}(r) \frac{d^2 \psi(r)}{d^2 r} + U(r) \psi(r) = -\Gamma \psi(r)$$

$$\psi(r) e^{2\Gamma t}$$

“magnetic energy density”

$$\kappa_{\text{diff}}(r) = \kappa_{\text{diff}}(\langle \delta v_i(\mathbf{r}_1, t) \delta v_j(\mathbf{r}_2, s) \rangle, \eta)$$

“mass”

$$U(r) = U(\langle \delta v_i(\mathbf{r}_1, t) \delta v_j(\mathbf{r}_2, s) \rangle, \eta)$$

“potential”

- can be solved with WKB approximation for large and small magnetic Prandtl numbers $P_m = R_m / Re$

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Critical Mag. Reynolds Number

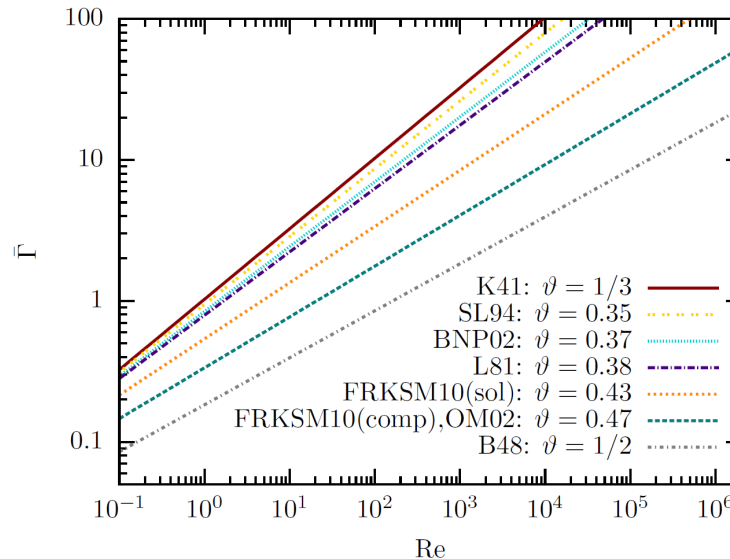
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- critical magnetic Reynolds number for small-scale dynamo action:
set $\Gamma = 0$ in Kazantsev equation and solve for R_m
($R_m = VL/\eta$)
- for Kolmogorov turbulence:
 $R_m > 110$
- for Burgers turbulence:
 $R_m > 2700$
- note: need high resolution in order to see dynamo in simulations

Growth Rate

- growth rate for large P_m :

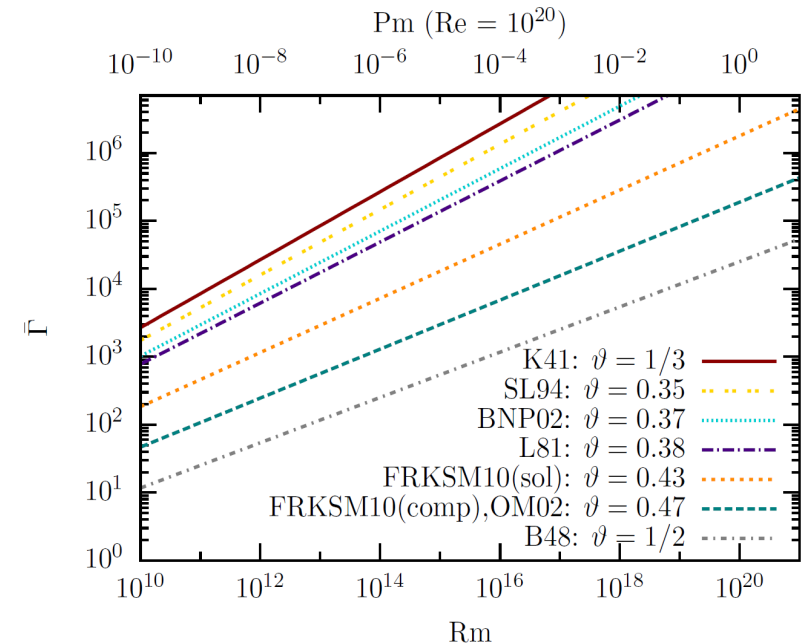
$$\Gamma \propto Re^{(1-\vartheta)/(1+\vartheta)}$$



Schober et al., PRE, 2012

- growth rate for small P_m :

$$\Gamma \propto Rm^{(1-\vartheta)/(1+\vartheta)}$$

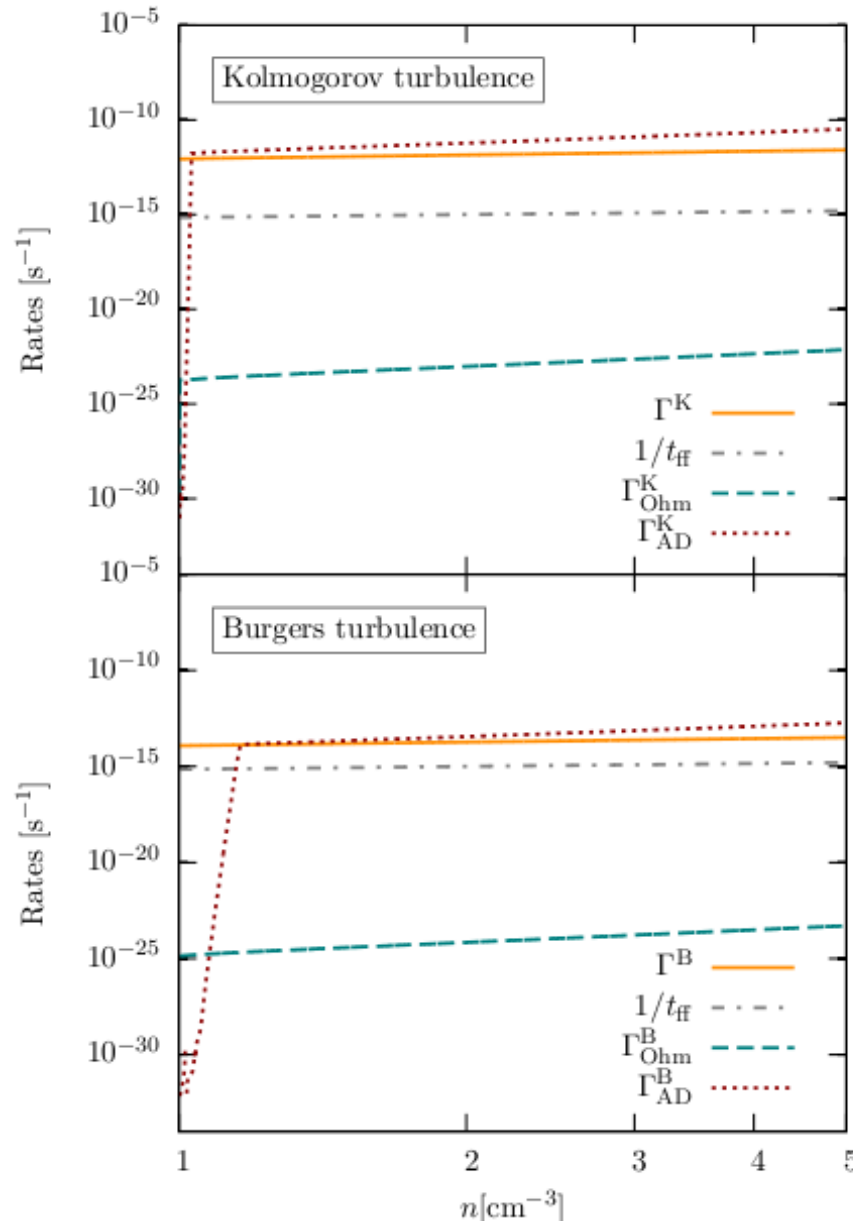


Schober et al., submitted

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Different Rates on Viscous Scale

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- growth rate:

$$\Gamma^K \propto \text{Re}^{1/2}$$

$$\Gamma^B \propto \text{Re}^{1/3}$$

- diffusion rates:

$$\Gamma_{\text{Ohm}} \approx \frac{\eta_{\text{Ohm}}}{\ell^2}$$

$$\Gamma_{\text{AD}} \approx \frac{\eta_{\text{AD}}}{\ell^2}$$

- inverse free-fall time:

$$\frac{1}{t_{\text{ff}}} = \left(\frac{3\pi}{32Gmn} \right)^{-1/2}$$

Magnetic Energy on Viscous Scale

- **small-scale magnetic energy:**

$$\frac{dE_B}{dt} = \left[\Gamma + \frac{4}{3n} \frac{dn}{dt} - \Gamma_{\text{Ohm}} - \Gamma_{\text{AD}}(E_B) \right] E_B$$

growth rate
of the small-
scale dynamo

spherical
gravitational
compression

$$B \propto n^{2/3}$$

Ohmic
dissipation

$$\Gamma_{\text{Ohm}} \approx \frac{\eta_{\text{Ohm}}}{\ell^2}$$

ambipolar
diffusion

$$\Gamma_{\text{AD}} \approx \frac{\eta_{\text{AD}}}{\ell^2}$$

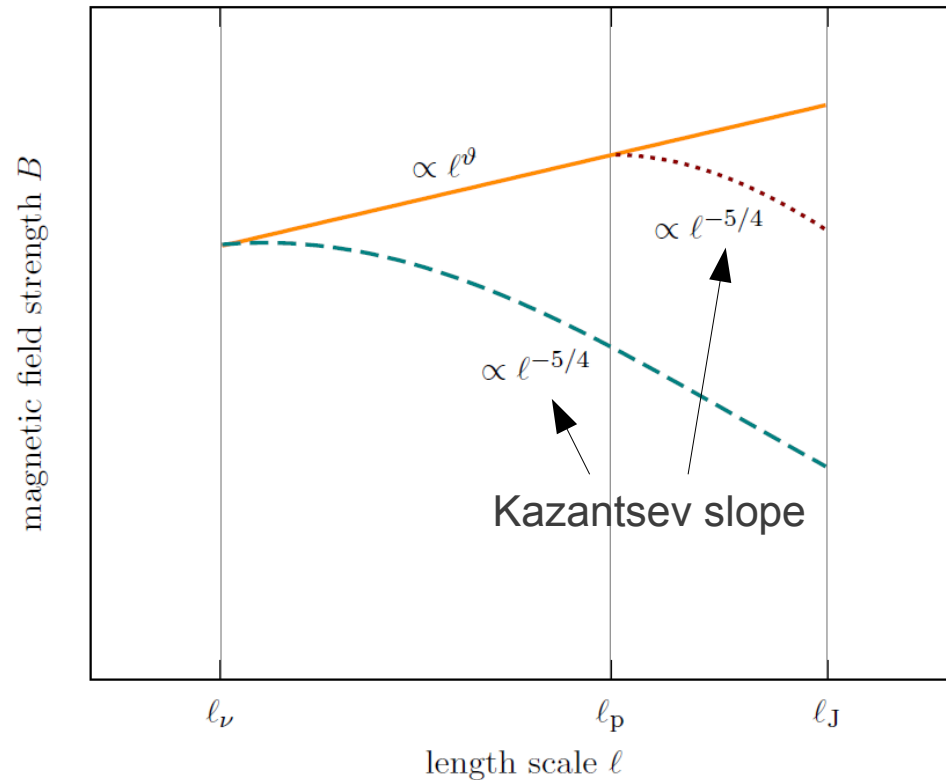
- **magnetic field strength:**

$$B = (8\pi E_B)^{1/2}$$

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Transport to Larger Scales: Non-Linear Dynamo

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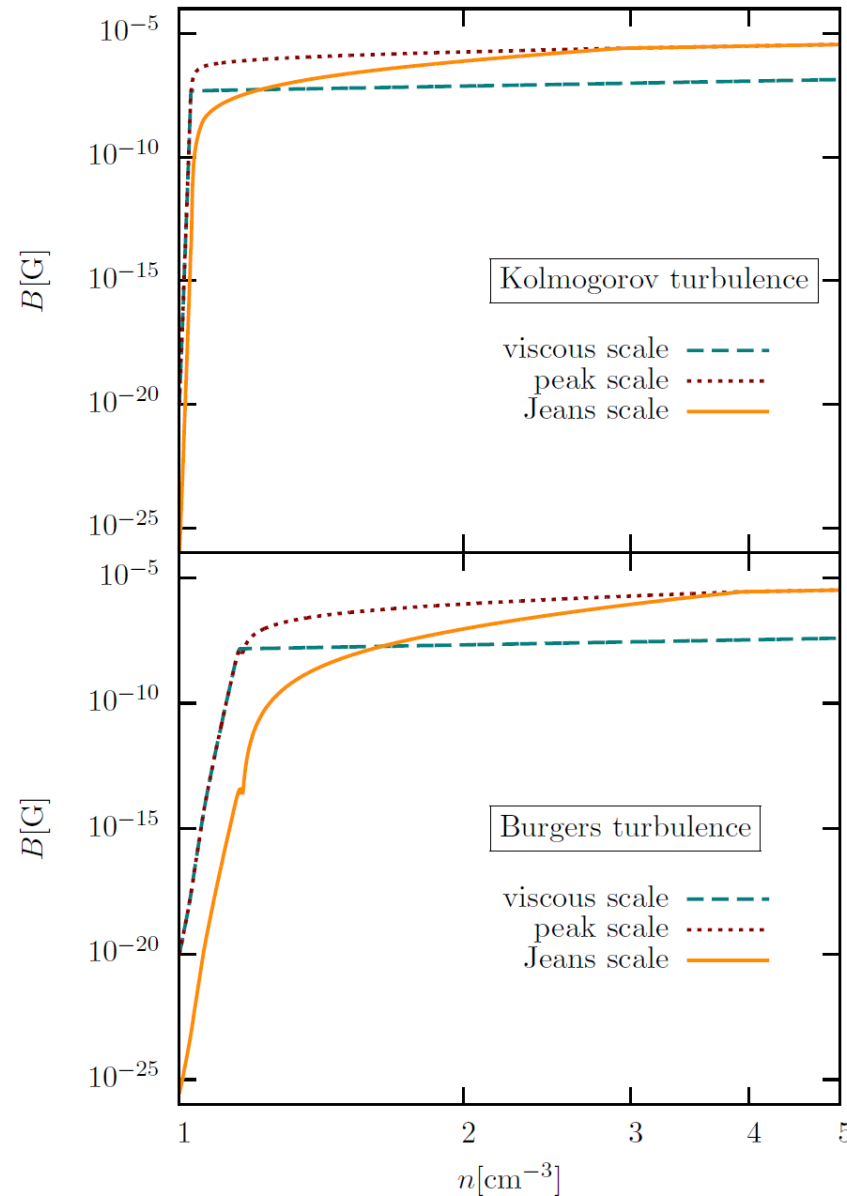
- after saturation on viscous (small) scale peak moves to larger scales
- assumption: peak moves to larger scales on the eddy timescale [Schekochihin et al., New J. Phys., 2002]
- peak scale:

$$l_p(t) = l_\nu(t_\nu) + \left(\frac{v_J}{\ell_J^\vartheta} (t - t_\nu) \right)^{1/(1-\vartheta)}$$

=> calculation of magnetic energy on the Jeans scale

Magnetic Field Evolution

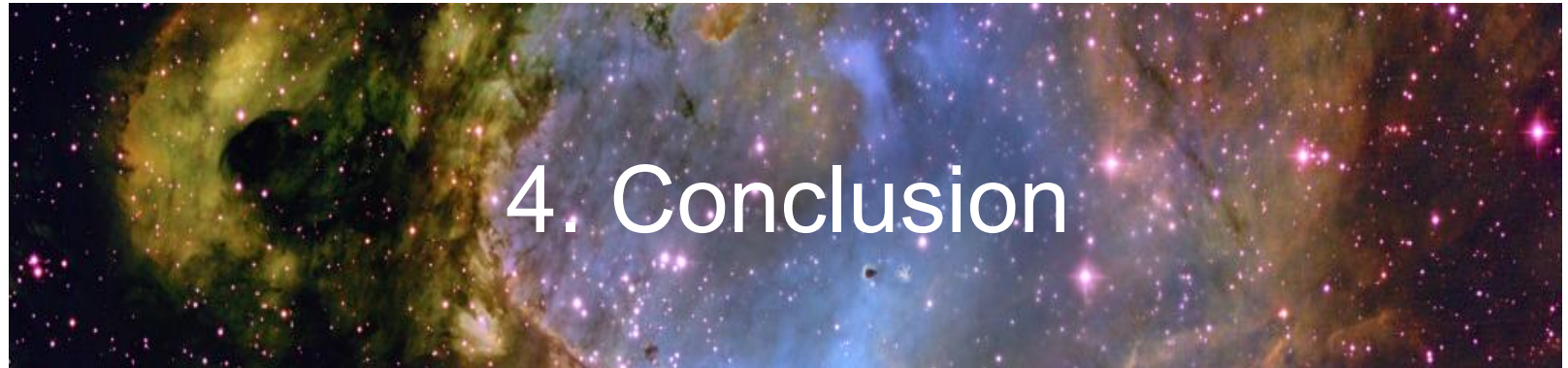
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=> Magnetic field saturates almost instantly (in terms of density) for all types of turbulence !!!
[Schober et al., ApJ, 2012]

“Small-Scale Dynamo & First Stars”

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- **critical magnetic Reynolds number** for small-scale dynamo action:

$$R_m > 110$$

Kolmogorov turbulence ($\vartheta = 1/3$)

$$R_m > 2700$$

Burgers turbulence ($\vartheta = 1/2$)

- **growth rate** of magnetic energy:

$$\Gamma \propto Re^{(1-\vartheta)/(1+\vartheta)}$$

large magnetic Prandtl number

$$\Gamma \propto R_m^{(1-\vartheta)/(1+\vartheta)}$$

small magnetic Prandtl number

- small-scale dynamo in **primordial star formation**:

Dynamical important magnetic fields can be generated on small timescales compared to the free-fall time.



Thanks for your attention!

Further questions/comments? →
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